If we consider infinite-range intermolecular potentials, then we have fractional powers rather than logarithms.

All the considerations of this section have the important consequence that approximate methods of solution which are not able to allow for a nonanalytic behavior for Kn $\to \infty$ produce poor results for large Knudsen numbers.

14. Perturbations of equilibria

The first steady solutions other than Maxwellian to be investigated were perturbations of the latter. The method of perturbation of equilibria is different from the Hilbert method because the small parameter is not contained in the Boltzmann equation but in auxiliary conditions, such as boundary or initial conditions. The advantage of the method is that we can investigate problems in the transition regime, provided differences in temperature and speed are moderate.

Let us try to find a solution of our problem for the Boltzmann equation in the form

$$f = \sum_{n=0}^{\infty} \varepsilon^n f_n, \tag{14.1}$$

where at variance with previous expansions ε is a parameter which *does not* appear in the Boltzmann equation. In addition f_0 is assumed from the start to be a Maxwellian distribution.

By inserting this formal series into Equation (5.1) and matching the various orders in ε, we obtain equations which one can hope to solve recursively:

$$\partial_t f_1 + \boldsymbol{\xi} \cdot \partial_{\mathbf{x}} f_1 = 2Q(f_1, f_0), \tag{14.2}_1$$

$$\ldots\ldots$$

$$\partial_t f_j + \boldsymbol{\xi} \cdot \partial_{\mathbf{x}} f_j = 2Q(f_j, f_0) + \sum_{i=1}^{j-1} Q(f_i, f_{j-i}), \tag{14.2}_j$$

$$\ldots\ldots$$

where, as in Section 4, $Q(f, g)$ denotes the symmetrized collision operator and the sum is empty for $j = 1$.

Although in principle one can solve the subsequent equations by recursion, in practice one solves only the first equation, which is called the *linearized Boltzmann equation*. This equation can be rewritten as follows:

$$\partial_t h + \boldsymbol{\xi} \cdot \partial_{\mathbf{x}} h = L_M h, \tag{14.3}$$

where L_M denotes the linearized collision operator about the Maxwellian M, i.e., $L_M h = 2Q(M, Mh)/M$, $h = f_1/M$ (see Section 10). We shall assume, as is usually done with

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Peipei Pang

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